

R-TORSION AND ZETA FUNCTIONS FOR ANALYTIC ANOSOV FLOWS ON 3-MANIFOLDS

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ABSTRACT. We improve previous results relating R-torsion, for an acyclic representation of the fundamental group, with a special value of the torsion zeta function of an analytic Anosov flow on a 3-manifold. By using the new techniques of Rugh and Fried we get rid of the unpleasant assumptions about the regularity of the invariant foliations.

1. INTRODUCTION

Let ϕ be a transitive analytic Anosov flow on a closed 3-manifold M . A representation $\rho : \pi_1(M) \rightarrow U(n)$ defines a flat vector bundle E_ρ . The torsion function is defined as the Euler product over the prime periodic orbits γ of period $l(\gamma)$

$$R_{\phi,\rho}(z) = \prod_{\gamma} \det(I - \Delta(\gamma)e^{-zl(\gamma)}\rho(\gamma))^{-1}$$

where $\Delta(\gamma) = \pm 1$ according to whether or not $E^u(\gamma)$ is orientable; in other words Δ is the holonomy of the orientation bundle $|E^u|$ of the unstable bundle. The product converges for $\Re z \gg 0$ and extending ideas of Rugh [7], [8], Fried [3] has constructed a meromorphic continuation to the whole complex plane.

Consider now a cell decomposition of M and the corresponding cell decomposition of the universal covering \tilde{M} . Let C^i be the set of equivariant \mathbb{C}^n valued i-cochains of \tilde{M} . Since equivariance is preserved by the coboundary map we obtain a cochain complex

$$C^0 \xrightarrow{d} C^1 \xrightarrow{d} C^2 \xrightarrow{d} C^3.$$

If its cohomology $H^*(M, E_\rho)$ vanishes, the representation is called *acyclic*. In such a case there is a chain retraction $D : C^{i+1} \rightarrow C^i$, then an isomorphism $d + D : C^0 \oplus C^2 \rightarrow C^1 \oplus C^3$ and using a geometrical basis one defines *R-torsion* as

$$\tau_\rho(M) = |\det(d + D)|.$$

Extending previous results [1], [9], we prove in this paper

Theorem 1. *Let ϕ be an analytic transitive Anosov flow on an orientable closed 3-manifold M . Let $\rho : \pi_1(M) \rightarrow U(n)$ be an acyclic representation. Suppose there is*

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a periodic orbit γ_0 such that $1, \Delta(\gamma_0)$ are not eigenvalues of $\rho(\gamma_0)$. Then the torsion function $R_{\phi, \rho}$ is regular at the origin and

$$\tau_\rho(M) = |R_{\phi, \rho}(0)|.$$

The typical example where this result applies is the unit tangent bundle $M = US$ of an orientable closed surface S with a C^ω Riemannian metric of negative curvature and ϕ the geodesic flow on M .

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2. R-TORSION

Along this paper ϕ will be a transitive analytic Anosov flow on an orientable closed 3-manifold M so that the tangent bundle has a continuous splitting $TM = E^u \oplus E^s \oplus E^c$ where E^u, E^s are flow invariant line bundles exponentially contracted by $D\phi_{-t}$ and $D\phi_t$ respectively as $t \rightarrow \infty$ and E^c is the line bundle generated by the vector field defining the flow.

Given a periodic orbit γ and $x \in \gamma$, we consider the derivative P_γ of the Poincaré map defined on $E_x^s \oplus E_x^u$, which can be represented in matrix form as $\begin{pmatrix} S_\gamma & 0 \\ 0 & U_\gamma \end{pmatrix}$. Defining

$$\Phi_s(x) = \lim_{t \rightarrow 0} \frac{1}{t} \log |D\phi_t|E_x^s| \quad \text{and} \quad \Phi_u(x) = \lim_{t \rightarrow 0} \frac{1}{t} \log |D\phi_t|E_x^u|,$$

we have

$$\log |S_\gamma| = \int_0^{l(\gamma)} \Phi_s(\phi_t(x)) dt \quad \text{and} \quad \log |U_\gamma| = \int_0^{l(\gamma)} \Phi_u(\phi_t(x)) dt.$$

We now introduce symbolic dynamics. According to Ratner [5] given $\eta > 0$ there is a Markov family of size η , i.e. a finite collection of compact sets $\{R_k\}_{k \in I}$ contained in a family of analytic transverse discs $\{D_j\}$ with $\text{diam}(R_k) < \eta$ such that

- 1) $\mathcal{R} = \cup_{k \in I} R_k$ meets every flow line in time $< \eta$.
- 2) There is a transition set $A \subset I \times I$ such that for each $(k, l) \in A$ there is a map $f_{kl} : R_{kl} \subset R_k \rightarrow R_l$ given by $f_{kl}(x) = \phi_{t_{kl}(x)}(x)$ for $x \in \text{int}R_{kl}$, where $t_{kl} : \text{int}R_{kl} \rightarrow]0, \eta[$ is the first return time to \mathcal{R} that extends to a continuous function on R_{kl} . The families $\{\text{int}R_{kl}\}$ and $\{\text{int}f(R_{kl})\}$ are both disjoint and

$$R_k = \bigcup_{l: (k, l) \in A} R_{kl}, \quad R_l = \bigcup_{k: (k, l) \in A} f_{kl}(R_{kl}).$$

Let $B_{kl} = \{\phi_t(x) : 0 \leq t < t_{kl}(x), x \in R_{kl}\}$ and $B_k = \bigcup_{l: (k, l) \in A} B_{kl}$. Then $M = \bigcup_{k \in I} B_k$.

Let γ_0 be an orbit as in Theorem 1. Ratner constructed Markov families $\{R_{kl}\}_{k \in I}$ of small size such that for any $k \in I, \partial R_k \subset W^s(\gamma_0) \cup W^u(\gamma_0)$. We will always use such Markov families $\{R_k\}$ of sufficiently small size to have each B_k simply connected. Fix a point $q \in R_J$ and for each $k \in I, k \neq J$ choose a point $p_k \in R_k$. For $(k, l) \in A$ choose a segment α_{kl} in B_{kl} joining p_k and p_l , then the union of these segments forms a direct graph which is connected since the flow is transitive. A maximal tree based at q gives a directed path α_k from q to p_k . For $(k, l) \in A$ let β_{kl} be the class of $\alpha_l^{-1} \alpha_{kl} \alpha_k$ in $\pi_1(M, q)$. Let $\rho : \pi_1(M) \rightarrow U(n)$ be an acyclic

representation, $\Delta : \pi_1(M) \rightarrow U(1)$ be the holonomy of the orientation bundle of the unstable bundle. Put $\rho_{kl} = \rho(\beta_{kl}), \Delta_{kl} = \Delta(\beta_{kl})$. Let

$$A = W^s(\gamma_0) \cap \bigcup_{k \in I} \partial B_k.$$

For $\epsilon > 0$ sufficiently small the interval bundle

$$W(\epsilon) := \bigcup_{x \in A} W^s(x; \epsilon)$$

has no self-intersections and then it is a tubular neighborhood of γ_0 . Hence, the manifold with piecewise smooth boundary $M(\epsilon) := M \setminus W(\epsilon)$ can be taken as the exterior of γ_0 . Since 1 is not an eigenvalue of $\rho(\gamma_0)$, the excision property of R-torsion [2] implies that $E_\rho(M, M(\epsilon))$ is acyclic with R-torsion given by

$$\tau_\rho(M, M(\epsilon)) = \tau_\rho(\overline{W(\epsilon)}, \partial \overline{W(\epsilon)}) = |\det(I - \rho(\gamma_0))|.$$

Since ρ is acyclic, so is $E_\rho(M(\epsilon))$ and we have

$$(1) \quad \tau_\rho(M) = |\det(I - \rho(\gamma_0))| \tau_\rho(M(\epsilon)).$$

Collapsing $M(\epsilon)$ along the leaves of the stable foliation we get a 2-dimensional complex Ω simply homotopic to $M(\epsilon)$ and therefore $E_\rho|\Omega$ is acyclic with

$$\tau_\rho(M(\epsilon)) = \tau_\rho(\Omega).$$

Each $B_k \cap M(\epsilon)$ collapses to a 2-cell T_k , the rectangles $R_k \cap M(\epsilon)$ collapse to unstable intervals U_k , and the intersection of $M(\epsilon)$ with the union of the stable walls of the B_k 's collapses to a directed graph $\Gamma \subset \Omega$ with set of vertices $J =$ boundary points of the U_k 's. For $p \in J$, denote by L_p the segment of Γ from p to the next point $g(p) \in J$ in the direction of the flow. Thus, periodic orbits of g are in 1 - 1 correspondence with the loops in Γ . We have the chain complex

$$0 \rightarrow \bigoplus_{k \in I} \mathbb{Z}T_k \xrightarrow{\partial_1} \bigoplus_{k \in I} \mathbb{Z}U_k \oplus \bigoplus_{p \in J} \mathbb{Z}L_p \xrightarrow{\partial_0} \bigoplus_{p \in J} \mathbb{Z}p \rightarrow 0$$

with

$$\partial_1(T_k) = \left(\sum_{l:(k,l) \in A} \Delta_{kl} U_l - U_k, h(T_k) \right),$$

$$\partial_0(L_p) = g(p) - p.$$

Putting $V = \mathbb{C}^n$, the corresponding twisted cochain complex is given by

$$(2) \quad V^J \xrightarrow{d_0} V^I \oplus V^J \xrightarrow{d_1} V^I$$

$$d_0 = \begin{pmatrix} b \\ G - I \end{pmatrix}, \quad d_1 = (H - I, h^*)$$

b, h^* are maps we don't need to know explicitly, $(Gx)_p = \rho(L_p)^t x_{g(p)}$ and

$$(Hy)_k = \sum_{l:(k,l) \in A} \Delta_{kl} \rho_{kl}^t y_l.$$

In [9] we proved the following proposition.

Proposition 1. *$G - I$ and $H - I$ are nonsingular and*

$$\tau_\rho(M) = \frac{|\det(I - \rho(\gamma_0)) \det(I - G)|}{|\det(I - H)|} = \frac{|\det(I - \rho(\gamma_0))|^2 |\det(I - \Delta(\gamma_0)\rho(\gamma_0))|}{|\det(I - H)|}.$$

3. ZETA FUNCTIONS

As in section 2, let $\{R_k\}$ be a Markov family with transition set A . Consider the usual subshift of finite type (Σ_A, σ) and the natural map $\Pi_A : \Sigma_A \rightarrow \mathcal{R}$ such that $\Pi_A(\mathbf{k})$ is the point whose trajectory runs through $\cdots R_{k_{-1}}; R_{k_0}, R_{k_1} \cdots$. Defining $\tau : \Sigma_A \rightarrow]0, \infty[$ by $\tau(\mathbf{k}) = t_{k_0 k_1}(\Pi_A(\mathbf{k}))$, one constructs the suspension flow (Σ, ψ) with first return time τ . Π_A determines a semiconjugacy Π between ψ and ϕ mapping periodic orbits of ψ to periodic orbits of ϕ . For $\sigma^m(\mathbf{k}) = \mathbf{k}$ let

$$\begin{aligned} \rho^m(\mathbf{k}) &= \rho_{k_{m-1}k_m} \cdots \rho_{k_0 k_1}, & \tau^m(\mathbf{k}) &= \tau(\mathbf{k}) + \cdots + \tau(\sigma^{m-1}\mathbf{k}), \\ \Delta^m(\mathbf{k}) &= \Delta_{k_0 k_1} \cdots \Delta_{k_{m-1} k_m}, \\ Df_{\mathbf{k}}^m &= Df_{k_{m-1}k_0} \cdots Df_{k_0 k_1}(\Pi_A(\mathbf{k})). \end{aligned}$$

For the suspension flow we have the torsion function

$$R_{\psi, \rho}(z) = \exp \sum_{m=1}^{\infty} \frac{1}{m} \sum_{\sigma^m \mathbf{k} = \mathbf{k}} \exp(-z\tau^m(\mathbf{k})) \Delta^m(\mathbf{k}) \text{tr} \rho^m(\mathbf{k}).$$

There is an overcounting of periodic orbits in the symbolic flow and to arrange for this one uses the auxiliary subshifts of Manning and Bowen. For our Markov families we only need to introduce a finite number of extra factors all coming from the original orbit γ_0 . As we saw in [9] one can write

$$(3) \quad R_{\phi, \rho}(z) = R_{\psi, \rho}(z) \det(I - e^{-zI(\gamma_0)} \Delta(\gamma_0) \rho(\gamma_0)) \det(I - e^{-zI(\gamma_0)} \rho(\gamma_0))^2.$$

For future reference we introduce the twisted Rugh zeta function

$$d_{\rho}(s) = \exp - \sum_{m=1}^{\infty} \frac{s^m}{m} \sum_{\sigma^m \mathbf{k} = \mathbf{k}} \frac{\text{tr} \rho^m(\mathbf{k})}{|\det(I - Df_{\mathbf{k}}^m)|}.$$

Recall that if $g : \Sigma_A \rightarrow \mathbb{R}$ is a Hölder function we define the zeta function

$$\zeta(s, g) = \exp \sum_{m=1}^{\infty} \frac{s^m}{m} \sum_{\sigma^m \mathbf{k} = \mathbf{k}} \exp(g^m(\mathbf{k}))$$

where $g^m(\mathbf{k}) = g(\mathbf{k}) + \cdots + g(\sigma^{m-1}\mathbf{k})$. We can choose [4] a Hölder function $G : \Sigma_A \rightarrow \mathbb{R}$ cohomologous to g and depending only on future coordinates. Since $g^m(\mathbf{k}) = G^m(\mathbf{k})$ whenever $\sigma^m \mathbf{k} = \mathbf{k}$ we have that $\zeta(s, g) = \zeta(s, G)$ and so we can think that g depends only on future coordinates. The radius of convergence of $\log \zeta(s, g)$ is $\exp(-P(g))$ where $P(g)$ is the topological pressure of g .

For $\mathbf{k} \in \Sigma_A$ and $i = u, s$ let

$$\lambda_i(\mathbf{k}) = \int_0^{\tau(\mathbf{k})} \Phi_i \circ \phi_t \circ \Pi_A(\mathbf{k}) dt,$$

then λ_u, λ_s are Hölder functions, and for $\sigma^m(\mathbf{k}) = \mathbf{k}$ we have

$$\det(I - Df_{\mathbf{k}}^m) = (1 - \lambda_u^m(\mathbf{k}))(1 - \lambda_s^m(\mathbf{k})).$$

To study the convergence of $d_{\rho}(s)$ we introduce the following functions

$$\begin{aligned} Z_1(s) &= \exp \sum_{m=1}^{\infty} \frac{s^m}{m} \sum_{\sigma^m \mathbf{k} = \mathbf{k}} \frac{\exp(-\lambda_u^m(\mathbf{k})) \text{tr} \rho^m(\mathbf{k})}{(\Delta^m(\mathbf{k}) \exp \lambda_u^m(\mathbf{k}) - 1)}, \\ Z_2(s) &= \exp \sum_{m=1}^{\infty} \frac{s^m}{m} \sum_{\sigma^m \mathbf{k} = \mathbf{k}} \frac{\Delta^m(\mathbf{k}) \exp(\lambda_s^m(\mathbf{k})) \text{tr} \rho^m(\mathbf{k})}{|\det(I - Df_{\mathbf{k}}^m)|}, \end{aligned}$$

$$Z(s, \rho) = \exp \sum_{m=1}^{\infty} \frac{s^m}{m} \sum_{\sigma^m \mathbf{k}=\mathbf{k}} \exp(-\lambda_u^m(\mathbf{k})) \text{tr} \rho^m(\mathbf{k}).$$

We have $1/d_\rho(s) = Z_1(s)Z_2(s)Z(s, \rho)$ and

$$\begin{aligned} & \lim_{m \rightarrow \infty} \frac{1}{m} \log \sum_{\sigma^m \mathbf{k}=\mathbf{k}} \frac{\exp(-\lambda_u^m(\mathbf{k}))}{|\Delta^m(\mathbf{k}) \exp \lambda_u^m(\mathbf{k}) - 1|} \\ &= \lim_{m \rightarrow \infty} \frac{1}{m} \log \sum_{\sigma^m \mathbf{k}=\mathbf{k}} \exp(-2\lambda_u^m(\mathbf{k})) = P(-2\lambda_u) < 0, \\ & \lim_{m \rightarrow \infty} \frac{1}{m} \log \sum_{\sigma^m \mathbf{k}=\mathbf{k}} \frac{\exp(\lambda_s^m(\mathbf{k}))}{|\det(I - Df_{\mathbf{k}}^m)|} \\ &= \lim_{m \rightarrow \infty} \frac{1}{m} \log \sum_{\sigma^m \mathbf{k}=\mathbf{k}} \exp(\lambda_s^m(\mathbf{k}) - \lambda_u^m(\mathbf{k})) = P(\lambda_s - \lambda_u) < 0, \\ & \lim_{m \rightarrow \infty} \frac{1}{m} \log \sum_{\sigma^m \mathbf{k}=\mathbf{k}} \exp(-\lambda_u^m(\mathbf{k})) = P(-\lambda_u) = 0. \end{aligned}$$

Therefore $\log Z_1, \log Z_2$ have radius of convergence greater than 1. $\log Z(s, \rho)$ converges absolutely on the unit disc and so does $\log d_\rho$.

We now describe the techniques of Rugh [7], [8] and Fried [3]. First we introduce hyperbolic correspondences. We say that a correspondence $f \subset (I \times J) \times (I' \times J')$ admits a *cross map* $c = (u, s) : I' \times J \rightarrow I \times J'$ if

$$f = \{(u(x', y), y, x', s(x', y)) : x' \in I', y \in J\}.$$

If s, u are restrictions of analytic contractions we say that f is an analytic hyperbolic correspondence. When $\partial_1 s = \partial_2 u = 0$, we say that f is *split*. We will consider compositions of nearly split correspondences as follows.

Proposition 2. [3] *Suppose that for $k = 0, 1, 2, \dots$ we have rectangles $I_k \times J_k$ and correspondences $f_k \subset (I_k \times J_k) \times (I_{k+1} \times J_{k+1})$ which admit cross maps $c_k = (u_k, s_k)$. Put*

$$\alpha_k = \|\partial_1 u_k\|, \beta_k = \|\partial_2 u_k\|, \gamma_k = \|\partial_1 s_k\|, \delta_k = \|\partial_2 s_k\|.$$

There is a positive function $\rho(\sigma)$ such that if $\alpha_k, \delta_k \leq \sigma < 1$ and $\beta_k, \gamma_k \leq \rho(\sigma)$ for all k then each composition $f_{m-1} \circ \dots \circ f_0$ is a hyperbolic correspondence with cross map $c_{(0, \dots, m-1)} = (u_{(0, \dots, m-1)}, s_{(0, \dots, m-1)})$.

If the size of a Markov family $\{R_k\}$ with $R_k \subset D_j(k)$ is sufficiently small, each correspondence f_{kl} extends to a hyperbolic correspondence between rectangles $I_k \times J_k$ and $I_l \times J_l$ with cross maps $c_{kl} = (u_{kl}, s_{kl})$ where I_k, J_k are intervals in $T_{p_k} D_{j(k)} \cap (E^c \oplus E^u)_{p_k}$ and $T_{p_k} D_{j(k)} \cap (E^c \oplus E^s)_{p_k}$ such that $I_k \times J_k \subset T_{p_k} D_{j(k)}$. To construct meromorphic continuation of the torsion function one considers sections of $|E^u| \otimes \wedge^i(TM/E^c) \otimes E_\rho|_{D_j}$. For $(k, l) \in A$ the transitions from B_l to B_k for $\wedge^i(TM/E^c)$ are given by

$$M_{kl}(1; x_l, y_k) = \frac{1}{\partial_1 u_{kl}} \begin{pmatrix} 1 & \partial_1 s_{kl} \\ -\partial_2 u_{kl} & \det Dc_{kl} \end{pmatrix},$$

$$M_{kl}(0; x_l, y_k) = 1, \quad M_{kl}(2; x_l, y_k) = \det M_{kl}(1; x_l, y_k) = \frac{\partial_2 s_{kl}}{\partial_1 u_{kl}}.$$

Denote $\tau_{kl}(x_l, y_k) = t_{kl}(u_{kl}(x_l, y_k), y_k)$. One chooses discs W_k, V_k in \mathbb{C} containing I_k, J_k and holomorphic extensions of c_{kl} and τ_{kl} to $W_l \times V_k$.

Denote by $\hat{\mathbb{C}}$ the Riemann sphere. Let \mathcal{B}_k be the space of holomorphic functions on $(\hat{\mathbb{C}} \setminus \overline{W}_k) \times V_k$ vanishing on $\{\infty\} \times V_k$ that extend continuously to the boundary. Define $G_{kl}^z : W_k^* \times V_k \times W_l \times V_l^* \rightarrow \mathbb{C}$ where W_k^*, V_l^* are neighborhoods of $\hat{\mathbb{C}} \setminus W_k$ and $\hat{\mathbb{C}} \setminus V_l$ by

$$G_{kl}^z(w_k, v_k, w_l, v_l) = \frac{\partial_1 u_{kl}(w_l, v_k) \exp(-z\tau_{kl}(w_l, v_k))}{(w_k - u_{kl}(w_l, v_k))(v_l - s_{kl}(w_l, v_k))}.$$

One defines transfer operators $\mathcal{L}_j(z); j = 0, 1, 2$ on $\mathcal{U}_0 = \mathcal{U}_2 = \bigoplus_{k \in I} \mathcal{B}_k \otimes \mathbb{C}^n$ and $\mathcal{U}_1 = \bigoplus_{k \in I} \mathcal{B}_k \otimes \mathbb{C}^2 \otimes \mathbb{C}^n$ by

$$(\mathcal{L}_j(z)\xi)_k = \sum_{l:(k,l) \in A} \Delta_{kl} L_{kl}(j; z) \otimes \rho_{kl}^t \xi_l$$

where

(4)
$$L_{kl}(j; z)\eta_l(w_k, v_k) = \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} G_{kl}^z(w_k, v_k, w_l, v_l) M_{kl}(j; w_l, v_k) \eta_l(w_l, v_l)$$

Remark 1. Since $c_{kl} = (u_{kl}, s_{kl}), (k, l) \in A$ is a contraction, transfer operators improve the region of analyticity.

We have power series expansions for the "kernels" $G_{kl}^z M_{kl}$, for example

$$G_{kl}^0(w_k, v_k, w_l, v_l) M_{kl}(2; w_l, v_k) = \sum_{m_1, m_2=0}^{\infty} E_{m_1 m_2}^k(w_k, v_k) g_{m_1 m_2}^l(w_l, v_l)$$

where $E_{m_1 m_2}^k(w_k, v_k) = (w_k - a_k)^{-m_1 - 1} (v_k - b_k)^{m_2}$ with a_k, b_k the centers of W_k, V_k and

$$g_{m_1 m_2}^l(w_l, v_l) = \int_{\partial V_l} \frac{dy}{2\pi i} \frac{\partial_2 s_{kl}(w_l, y) (u_{kl}(w_l, y) - a_k)^{m_1}}{(v_l - s_{kl}(w_l, y))(y - b_k)^{m_2 + 1}}.$$

One sees [7] that $g_{m_1 m_2}^l$ are holomorphic on a neighborhood of $\overline{W}_l \times \hat{\mathbb{C}} \setminus V_l$ and that there is $r < 1$ such that $\|g_{m_1 m_2}^l\| < Cr^{m_1 + m_2}$. Thus we have convergent expansions

$$L_{kl}(2; 0)\eta_l = \sum_{m_1, m_2=0}^{\infty} E_{m_1 m_2}^k \langle g_{m_1 m_2}^l, \eta_l \rangle$$

where

$$\langle g_{m_1 m_2}^l, \eta_l \rangle = \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} g_{m_1 m_2}^l(w_l, v_l) \eta_l(w_l, v_l).$$

This implies that $\mathcal{L}_2(0)$ is a nuclear operator of order zero and therefore a trace class operator and we can form the Fredholm determinant

$$\det(I - \mathcal{L}_2(0)) = \exp - \sum_{m=1}^{\infty} \frac{\text{tr}(\mathcal{L}_2(0)^m)}{m} = \prod_{\lambda} (1 - \lambda)$$

where the product runs over the eigenvalues of $\mathcal{L}_2(0)$ counted with multiplicity.

If $\beta \in \mathcal{U}_2^*, \xi \in \mathcal{U}_2$ let $\beta|_{\mathcal{B}_k} \otimes \mathbb{C}^n = \sum_i \beta_k^i \otimes e_i$ and $\xi_k = \sum_j \xi_k^j \otimes e_j$ then

$$\langle \mathcal{L}_2(0)^* \beta, \xi \rangle = \sum_{(k,l) \in A} \sum_{m_1, m_2=0}^{\infty} \sum_{i,j=1}^n \Delta_{kl} \langle \beta_k^i, E_{m_1 m_2}^k \rangle \langle g_{m_1 m_2}^l, \xi_l^j \rangle \langle e_i, \rho_{kl}^t e_j \rangle.$$

Thus

$$(\mathcal{L}_2(0)^*\beta)|_{\mathcal{B}_l} \otimes \mathbb{C}^n = \sum_{k:(k,l) \in A} \sum_{m_1, m_2=0}^{\infty} \sum_{i=1}^n \Delta_{kl} \langle \beta_k^i, E_{m_1 m_2}^k \rangle g_{m_1 m_2}^l \otimes \overline{\rho_{kl}} e_i$$

is analytic on a neighborhood of $\overline{W}_l \times \hat{\mathbb{C}} \setminus V_l$.

Remark 2. Transfer operators have the functorial property.

This implies for $\mathbf{k} = (k_0, \dots, k_m)$ with $(k_i, k_{i+1}) \in A$

$$\begin{aligned} L_{\mathbf{k}}^m(j; z) \eta_{k_m}(w, v) &:= L_{k_0 k_1}(j; z) \cdots L_{k_{m-1} k_m}(j; z) \eta_{k_m}(w, v) \\ &= \int_{\partial W_{k_m}} \int_{\partial V_{k_m}} \frac{dw' dv'}{(2\pi i)^2} G_{\mathbf{k}}^z(w, v, w', v') M_{\mathbf{k}}(j; w', v) \eta_l(w', v') \end{aligned}$$

where $M_{\mathbf{k}}$ is defined as M_{kl} with c_{kl} replaced by $c_{\mathbf{k}} = (u_{\mathbf{k}}, s_{\mathbf{k}})$,

$$G_{\mathbf{k}}^z(w, v, w', v') = \frac{\partial_1 u_{\mathbf{k}}(w', v) \exp(-z\tau_{\mathbf{k}}(w', v))}{(w - u_{\mathbf{k}}(w', v))(v' - s_{\mathbf{k}}(w', v))},$$

and $\tau_{\mathbf{k}}$ is a holomorphic extension of

$$\tau_{\mathbf{k}}(x, y) = t_{k_0 k_1}(u_{\mathbf{k}}(x, y), y) + \cdots + t_{k_{m-1} k_m} \circ f_{k_{m-2} k_{m-1}} \circ \cdots \circ f_{k_0 k_1}(u_{\mathbf{k}}(x, y), y).$$

We have

$$\text{tr}(\mathcal{L}_j(z)^m) = \sum_{\sigma^m \mathbf{k} = \mathbf{k}} \Delta^m(\mathbf{k}) \text{tr} \rho^m(\mathbf{k}) \text{tr} L_{\mathbf{k}}^m(j; z).$$

As proved by Rugh [7] and Fried [3] for $\sigma^m \mathbf{k} = \mathbf{k}$ we have

$$\text{tr} L_{\mathbf{k}}^m(j; z) = - \frac{\exp(-z\tau^m(\mathbf{k})) \text{tr}(\wedge^j Df_{\mathbf{k}}^m)}{\det(I - Df_{\mathbf{k}}^m)}.$$

Thus

$$\text{tr}(\mathcal{L}_0(0)^m) = \sum_{\sigma^m \mathbf{k} = \mathbf{k}} \frac{\text{tr} \rho^m(\mathbf{k})}{|\det(I - Df_{\mathbf{k}}^m)|},$$

$$\text{tr}(\mathcal{L}_0(z)^m) + \text{tr}(\mathcal{L}_1(z)^m) + \text{tr}(\mathcal{L}_2(z)^m) = - \sum_{\sigma^m \mathbf{k} = \mathbf{k}} \exp(-z\tau^m(\mathbf{k})) \Delta^m(\mathbf{k}) \text{tr} \rho^m(\mathbf{k}).$$

Therefore the expression

$$(5) \quad R_{\psi, \rho}(z) = \frac{\det(I - \mathcal{L}_0(z)) \det(I - \mathcal{L}_2(z))}{\det(I - \mathcal{L}_1(z))}$$

provides a meromorphic continuation of $R_{\psi, \rho}(z)$. Similarly

$$d_{\rho}(s) = \det(I - s\mathcal{L}_0(0)) = \prod_{\nu} (1 - s\nu)$$

where the product is taken over the eigenvalues of $\mathcal{L}_0(0)$, gives an analytic continuation of $d_{\rho}(s)$.

4. SPECIAL VALUE OF THE TORSION FUNCTION

In this section we prove Theorem 1 by using (3) and (5) to evaluate $R_{\phi,\rho}(0)$. Proposition 3 says that this value is the same we obtained in [9] under the assumption that the stable foliation was analytic. To use (5) we need the following lemmas.

Lemma 1. *Let $\rho : \pi_1(M) \rightarrow U(n)$ be an acyclic representation. Then 1 is neither an eigenvalue of $\mathcal{L}_0(0)$ nor of $\mathcal{L}_2(0)$*

Proof. Suppose 1 is an eigenvalue of $\mathcal{L}_0(0)$, then $s = 1$ is a zero of $d_\rho(s)$ and so it is pole of $Z(s, \rho)$. We recall that in order to obtain a meromorphic continuation of $Z(s, \rho)$, one changes λ_u to a Hölder cohomologous function Λ_u , depending only on future coordinates, and considers the Ruelle operators \mathcal{N} and \mathcal{N}_ρ defined on spaces of Hölder functions on Σ_A^+ with values in \mathbb{C} and \mathbb{C}^n respectively by

$$(\mathcal{N}g)(\mathbf{k}) = \sum_{\sigma\mathbf{l}=\mathbf{k}} \exp(-\Lambda_u(\mathbf{l}))g(\mathbf{l})$$

$$(\mathcal{N}_\rho\eta)(\mathbf{k}) = \sum_{\sigma\mathbf{l}=\mathbf{k}} \exp(-\Lambda_u(\mathbf{l}))\rho_{l_0k_0}\eta(\mathbf{l}).$$

Since $P(-\Lambda_u) = 0$ we have that 1 is an eigenvalue of \mathcal{N} with a positive eigenfunction h . According to [4] and [6], $Z(s, \rho)$ has a meromorphic continuation to a disc of radius > 1 and its poles on that disc are the inverses λ^{-1} of the eigenvalues of \mathcal{N}_ρ ; therefore 1 is an eigenvalue of \mathcal{N}_ρ . Let ξ be an eigenvector and put $\eta = \xi/h$, then

$$(6) \quad \eta(\mathbf{k}) = \sum_{\sigma\mathbf{l}=\mathbf{k}} \frac{h(\mathbf{k})}{h(\mathbf{l})} \exp(-\Lambda_u(\mathbf{l}))\rho_{l_0k_0}\eta(\mathbf{l}).$$

Let $K = \max\{|\eta(\mathbf{l})| : \mathbf{l} \in \Sigma_A\}$ and $\mathbf{r} \in \Sigma_A$ such that $|\eta(\mathbf{r})| = K$. If $|\eta(\mathbf{l})| < K$ for some \mathbf{l} with $\sigma\mathbf{l} = \mathbf{r}$ then

$$|\eta(\mathbf{r})| \leq \sum_{\sigma\mathbf{l}=\mathbf{r}} \frac{h(\mathbf{r})}{h(\mathbf{l})} \exp(-\Lambda_u(\mathbf{l}))|\eta(\mathbf{l})| < K.$$

Therefore $|\eta(\mathbf{l})| = K$ if $\sigma^m\mathbf{l} = \mathbf{r}$. Since σ is transitive $|\eta(\mathbf{k})| = K$ for any $\mathbf{k} \in \Sigma_A$. Thus the left-hand side of (6) is a convex combination of points in the sphere $S = \{x \in \mathbb{C}^n : |x| = K\}$ and the right-hand side is a point in S . We then have $\rho_{l_0l_1}\eta(\mathbf{l}) = \eta \circ \sigma(\mathbf{l})$ for any $\mathbf{l} \in \Sigma_A$.

For $a \in I$ choose \mathbf{l} periodic, say $\sigma^m\mathbf{l} = \mathbf{l}$, such that $l_0 = a$ and let $\eta_a = \eta(\mathbf{l})$. If γ is the corresponding periodic orbit of the flow we have

$$\rho(\gamma)\eta_a = \eta_a.$$

Given $\mathbf{k} \in \Sigma_A$ with $k_0 = a$ define $\mathbf{k}^j = (k_0^j, k_1^j, \dots)$ by $k_s^j = l_s$ if $s < jm$ and $k_{s+jm}^j = k_s$. Then \mathbf{k}^j converges to \mathbf{l} and

$$\rho(\gamma)^j\eta(\mathbf{k}^j) = \eta(\mathbf{k}).$$

Since ρ is unitary

$$\lim_{j \rightarrow \infty} |\rho(\gamma)^j\eta(\mathbf{k}^j) - \eta_a| = \lim_{j \rightarrow \infty} |\rho(\gamma)^j(\eta(\mathbf{k}^j) - \eta_a)| = 0;$$

therefore $\eta(\mathbf{k}) = \eta_a$. Thus η defines a nontrivial element of $H^0(M, E_\rho)$, which contradicts the acyclicity of ρ .

Now suppose that 1 is an eigenvalue of the compact operator $\mathcal{L}_2(0)$ then 1 is also an eigenvalue of $\mathcal{L}_2(0)^*$ and let β be an eigenvector. According to Remark 1, β

belongs to the space $\mathcal{U}' = \bigoplus_{k \in I} \mathcal{B}'_k \otimes \mathbb{C}^n$ where \mathcal{B}'_k is the space of analytic functions on $W_k \times \hat{\mathbb{C}} \setminus \bar{V}_k$ that extend continuously to the boundary. Restricted to \mathcal{U}' , $\mathcal{L}_2(0)^*$ is given by

$$(\mathcal{L}_2(0)^*\zeta)_l = \sum_{k:(k,l) \in A} \Delta_{kl} L_{kl}^* \otimes \overline{\rho_{kl}} \zeta_k$$

where

$$L_{kl}^* \xi_l(w_l, v_l) = \int_{\partial W_k} \int_{\partial V_k} \frac{dw_k dv_k}{(2\pi i)^2} H_{kl}(w_k, v_k, w_l, v_l) \xi_l(w_k \cdot v_k),$$

$$H_{kl}(w_k, v_k, w_l, v_l) = \frac{\partial_2 s_{kl}(w_l, v_k)}{(w_k - u_{kl}(w_l, v_k))(v_l - s_{kl}(w_l, v_k))}.$$

We see that the restriction of $\mathcal{L}_2(0)^*$ to \mathcal{U}' is the operator $\mathcal{L}_0(0)$ for the reversed time flow and the contragradient representation $\rho^* = \rho$. As we have proved 1 is not an eigenvalue of this operator. \square

Lemma 2. *The holomorphic exterior derivative d intertwines transfer operators at $z = 0$ i.e. $\mathcal{L}_{j+1}(0)d = d\mathcal{L}_j(0)$ provided the left-hand side is defined.*

Proof. From (4) we have

$$L_{kl}(0; 0)\eta_l(w_k, v_k) = \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \frac{\partial_1 u_{kl}(w_l, v_k)\eta_l(w_l, v_l)}{(w_k - u_{kl}(w_l, v_k))(v_l - s_{kl}(w_l, v_k))},$$

$$\begin{pmatrix} \Xi_{kl}(w_k, v_k) \\ Z_{kl}(w_k, v_k) \end{pmatrix} := L_{kl}(1; 0) \begin{pmatrix} \xi_l \\ \zeta_l \end{pmatrix} (w_k, v_k)$$

$$= \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \frac{\begin{pmatrix} \xi_l + \partial_1 s_{kl} \zeta_l \\ -\xi_l \partial_2 u_{kl} + \det Dc_{kl} \zeta_l \end{pmatrix}}{(w_k - u_{kl}(w_l, v_k))(v_l - s_{kl}(w_l, v_k))},$$

$$L_{kl}(2; 0)\eta_l(w_k, v_k) = \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \frac{\partial_2 s_{kl}(w_l, v_k)\eta_l(w_l, v_l)}{(w_k - u_{kl}(w_l, v_k))(v_l - s_{kl}(w_l, v_k))}.$$

Integration by parts gives

$$\begin{aligned} \frac{\partial L_{kl}(0; 0)\eta_l}{\partial w_k} &= - \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \frac{\partial_1 u_{kl}(w_l, v_k)\eta_l(w_l, v_l)}{(w_k - u_{kl}(w_l, v_k))^2 (v_l - s_{kl}(w_l, v_k))} \\ &= \int_{\partial V_l} \frac{dv_l}{(2\pi i)^2} \int_{\partial W_l} \frac{dw_l}{w_k - u_{kl}(w_l, v_k)} \frac{\partial}{\partial w_l} \left(\frac{\eta_l(w_l, v_l)}{v_l - s_{kl}(w_l, v_k)} \right) \\ &= \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \left[\frac{\partial_1 \eta_l}{(w_k - u_{kl})(v_l - s_{kl})} + \frac{\eta_l \partial_1 s_{kl}}{(w_k - u_{kl})(v_l - s_{kl})^2} \right] \\ &= \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \frac{\partial_1 \eta_l + \partial_2 \eta_l \partial_1 s_{kl}}{(w_k - u_{kl})(v_l - s_{kl})}. \end{aligned}$$

$$\begin{aligned} \frac{\partial L_{kl}(0; 0)\eta_l}{\partial v_k} &= \int_{\partial W_l} \int_{\partial V_l} \left[\partial_2 \partial_1 u_{kl} + \frac{\partial_1 u_{kl} \partial_2 u_{kl}}{w_k - u_{kl}} + \frac{\partial_1 u_{kl} \partial_2 s_{kl}}{v_l - s_{kl}} \right] \\ &\quad \times \frac{\eta_l(w_l, v_l) dw_l dv_l}{(2\pi i)^2 (w_k - u_{kl})(v_l - s_{kl})} \\ &= \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \frac{-\partial_2 u_{kl} \partial_1 \eta_l + (\partial_1 u_{kl} \partial_2 s_{kl} - \partial_1 s_{kl} \partial_2 u_{kl}) \partial_2 \eta_l}{(w_k - u_{kl})(v_l - s_{kl})}. \end{aligned}$$

(8) and (9) imply $d\mathcal{L}_0(0) = \mathcal{L}_1(0)d$.

$$\begin{aligned} \frac{\partial \Xi_{kl}}{\partial v_k} &= \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \left[\frac{\partial_2 \partial_1 s_{kl} \zeta_l}{(w_k - u_{kl})(v_l - s_{kl})} \right. \\ &\quad \left. + (\xi_l + \partial_1 s_{kl} \zeta_l) \left(\frac{\partial_2 u_{kl}}{(w_k - u_{kl})^2 (v_l - s_{kl})} + \frac{\partial_2 s_{kl}}{(w_k - u_{kl})(v_l - s_{kl})^2} \right) \right], \\ \frac{\partial Z_{kl}}{\partial w_k} &= \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \frac{\xi_l \partial_2 u_{kl} - \det Dc_{kl} \zeta_l}{(w_k - u_{kl})^2 (v_l - s_{kl})} \end{aligned}$$

As before

$$\begin{aligned} (10) \quad \frac{\partial \Xi_{kl}}{\partial v_k} - \frac{\partial Z_{kl}}{\partial w_k} &= \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \left[\frac{\partial_1 u_{kl} \partial_2 s_{kl} \zeta_l}{(w_k - u_{kl})^2 (v_l - s_{kl})} \right. \\ &\quad \left. + \frac{(\xi_l + \partial_1 s_{kl} \zeta_l) \partial_2 s_{kl}}{(w_k - u_{kl})(v_l - s_{kl})^2} + \frac{\partial_2 \partial_1 s_{kl} \zeta_l}{(w_k - u_{kl})(v_l - s_{kl})} \right] \\ &= \int_{\partial W_l} \int_{\partial V_l} \frac{dw_l dv_l}{(2\pi i)^2} \frac{\partial_2 s_{kl}(w_l, v_k) (\partial_2 \xi_l(w_l, v_l) - \partial_1 \zeta_l(w_l, v_l))}{(w_k - u_{kl}(w_l, v_k))(v_l - s_{kl}(w_l, v_k))} \end{aligned}$$

which implies $d\mathcal{L}_1(0) = \mathcal{L}_2(0)d$. □

Propositions 1 and 3 (coming up) complete the proof of Theorem 1.

Proposition 3.

$$R_{\phi, \rho}(0) = \frac{\det(I - \Delta(\gamma_0)\rho(\gamma_0)) \det(I - \rho(\gamma_0))^2}{\det(I - H)}.$$

Proof. Let \mathcal{V}_j be the space of $\xi \in \mathcal{U}_j$ such that ξ_k is analytic in a neighborhood of $\hat{\mathbb{C}} \setminus W_k \times \bar{V}_k$.

Let $\lambda \neq 0$ be an eigenvalue of $\mathcal{L}_j(0)$. If $(\mathcal{L}_j(0) - \lambda I)\xi = 0$ then $\xi \in \mathcal{V}_j$. It follows by induction that $(\mathcal{L}_j(0) - \lambda I)^k \eta = 0$ implies that $\eta \in \mathcal{V}_j$ and $(\mathcal{L}_j(0) - \lambda I)^k d\eta = 0$.

If $(\mathcal{L}_2(0) - \lambda I)^k \eta = 0$ for $\lambda \neq 0$, there is $\xi \in \mathcal{V}_1$ such that $\eta = d\xi$. Therefore

$$\det(I - \mathcal{L}_1(0)) = \det(I - \mathcal{L}_2(0)) \det(I - \mathcal{L}_1(0)|_{\ker d \cap \mathcal{V}_1}).$$

Let

$$\omega_k = \frac{dw_k}{w_k - a_k} \quad \text{and} \quad \mathcal{H} = \{\xi \in \mathcal{U}_1 : \xi_k = c_k \omega_k, c_k \in \mathbb{C}^n\}.$$

It is straightforward to show that $\ker d \cap \mathcal{V}_1 = d\mathcal{V}_0 \oplus \mathcal{H}$ and for $\xi \in \mathcal{H}$ there is $\eta \in \mathcal{V}_0$ such that

$$(\mathcal{L}_1(0)\xi)_k = \sum_{l:(k,l) \in A} \Delta_{kl} \rho_{kl}^t c_l \omega_k + d\eta_k.$$

Since $\mathcal{L}_1(0)d = d\mathcal{L}_0(0)$ we have

$$\det(I - \mathcal{L}_1(0)|_{\ker d \cap \mathcal{V}_1}) = \det(I - \mathcal{L}_0(0)) \det(I - H)$$

and so

$$\det(I - \mathcal{L}_1(0)) = \det(I - \mathcal{L}_2(0)) \det(I - \mathcal{L}_0(0)) \det(I - H).$$

Since 1 is neither an eigenvalue of $\mathcal{L}_0(0)$ nor of $\mathcal{L}_2(0)$ and $I - H$ is nonsingular (Proposition 1) we have that 1 is not an eigenvalue of $\mathcal{L}_1(0)$. By (5) and (3) we finally obtain Proposition 3. □

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