

Non-commutative Painlevé equations and Hermite-type matrix orthogonal polynomials*

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*Joint work with Mattia Cafasso, Université d'Angers

OUTLINE

1. Motivation and preliminaries
2. Integral representations
3. Non-commutative Painlevé IV equation

MOTIVATION

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$$K_n(x, y) = \sum_{k=0}^{n-1} H_k(x) H_k(y) e^{-\frac{x^2+y^2}{2}}$$

describes the statistical properties of the eigenvalues of a **random matrix** \mathbf{M} in the space of $(n \times n)$ Hermitian matrices with the measure $\mu(\mathbf{M}) = e^{-\text{Tr}(\mathbf{M}^2)} d\mathbf{M}$ (GUE, Mehta).

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The **last particle distribution** is given by the **Fredholm determinant**

$$F(s) = \mathbb{P} [\lambda_{\max} \leq s] = \det(\text{Id} - \chi_s \mathbb{K}_n)$$

where χ_s is the indicator function of the interval $[s, \infty)$ and $\mathbb{K}_n : L^2(\mathbb{R}) \rightarrow L^2(\mathbb{R})$ is the **integral operator**

$$[\mathbb{K}_n f](x) = \int_{\mathbb{R}} K_n(x, y) f(y) dy \quad \forall f \in L^2(\mathbb{R})$$

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THEOREM (Tracy-Widom, 1994)

The log derivative of the Fredholm determinant

$$R(s) = \partial_s \log(\det(\text{Id} - \chi_s \mathbb{K}_n))$$

solves the **sigma-form of the Painlevé IV equation**

$$(R'')^2 + 4(R')^2(R' + 2n) - 4(sR' - R)^2 = 0$$

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Extend these results to CD kernels associated to Hermite-type **matrix-valued** orthogonal polynomials (MOP).

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 \Rightarrow **Matrix-valued CD kernels.**
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- **Non-commutative integrable systems**

P. Etingof, I. Gelfand, V. Retakh, V. Rubtsov, V. Sokolov...

PRELIMINARIES

Let \mathbf{W} be a **weight matrix** (positive definite and finite moments).

Consider $L^2_{\mathbf{W}}(\mathbb{R}, \mathbb{C}^{N \times N})$ the **weighted space** with the inner product

$$\langle \mathbf{F}, \mathbf{G} \rangle_{\mathbf{W}} = \int_{\mathbb{R}} \mathbf{F}(x) \mathbf{W}(x) \mathbf{G}^*(x) dx$$

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A sequence $(\mathbf{P}_n)_n$ of **matrix orthonormal polynomials (MOP)** with respect to \mathbf{W} is a sequence satisfying (Krein, 1949)

$$\deg \mathbf{P}_n = n, \quad \langle \mathbf{P}_n, \mathbf{P}_m \rangle_{\mathbf{W}} = \mathbf{I}_N \delta_{nm}$$



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If $(\mathbf{P}_n)_n$ is complete, the **Christoffel-Darboux (CD) kernel** is

$$\mathbf{K}_n(x, y) = \sum_{k=0}^{n-1} \mathbf{P}_k^*(y) \mathbf{P}_k(x), \quad x, y \in \mathbb{R}$$

SOME NOTATION

$$\mathbf{A}_N = \begin{pmatrix} 0 & \nu_1 & 0 & \cdots & 0 \\ 0 & 0 & \nu_2 & \cdots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \nu_{N-1} \\ 0 & 0 & 0 & \cdots & 0 \end{pmatrix}, \nu_i \in \mathbb{R}, \mathbf{J}_N = \begin{pmatrix} N-1 & \cdots & 0 & 0 \\ \vdots & \ddots & \vdots & \vdots \\ 0 & \cdots & 1 & 0 \\ 0 & \cdots & 0 & 0 \end{pmatrix}$$

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$$z^{\mathbf{J}} = \begin{pmatrix} z^{N-1} & & & \\ & \ddots & & \\ & & z & \\ & & & 1 \end{pmatrix}, \quad z^{-\mathbf{J}} = \begin{pmatrix} \frac{1}{z^{N-1}} & & & \\ & \ddots & & \\ & & \frac{1}{z} & \\ & & & 1 \end{pmatrix}$$

THE EXAMPLE

Let us consider the weight matrix (Durán-Grünbaum, 2004)

$$\mathbf{W}(x) = e^{-x^2} e^{\mathbf{A}x} e^{\mathbf{A}^*x}, \quad x \in \mathbb{R}$$

and the family of MOP $(\mathbf{P}_n)_n$
satisfying the second order differential equation

$$\mathbf{P}_n''(x) + \mathbf{P}_n'(x)(-2x\mathbf{I} + 2\mathbf{A}) + \mathbf{P}_n(x)(\mathbf{A}^2 - 2\mathbf{J}) = (-2n\mathbf{I} - 2\mathbf{J})\mathbf{P}_n(x)$$

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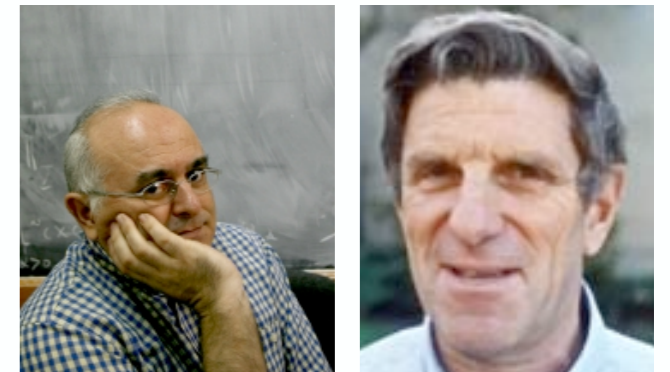


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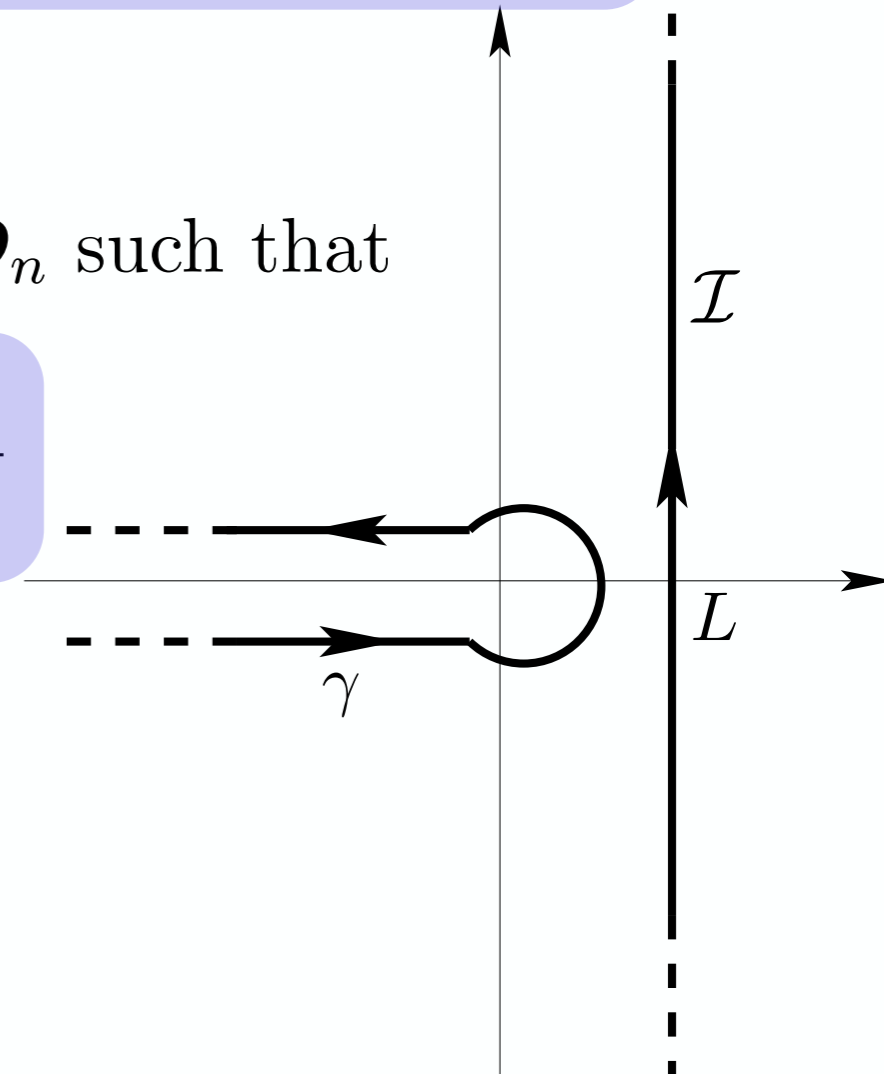


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THEOREM (Cafasso-MdI, 2013)

There exist suitable constant matrices C_n and D_n such that

$$P_n(x)e^{Ax} = \oint_{\gamma} z^{-J} C_n z^J e^{-z^2 + 2zx} \frac{dz}{z^{n+1}}$$



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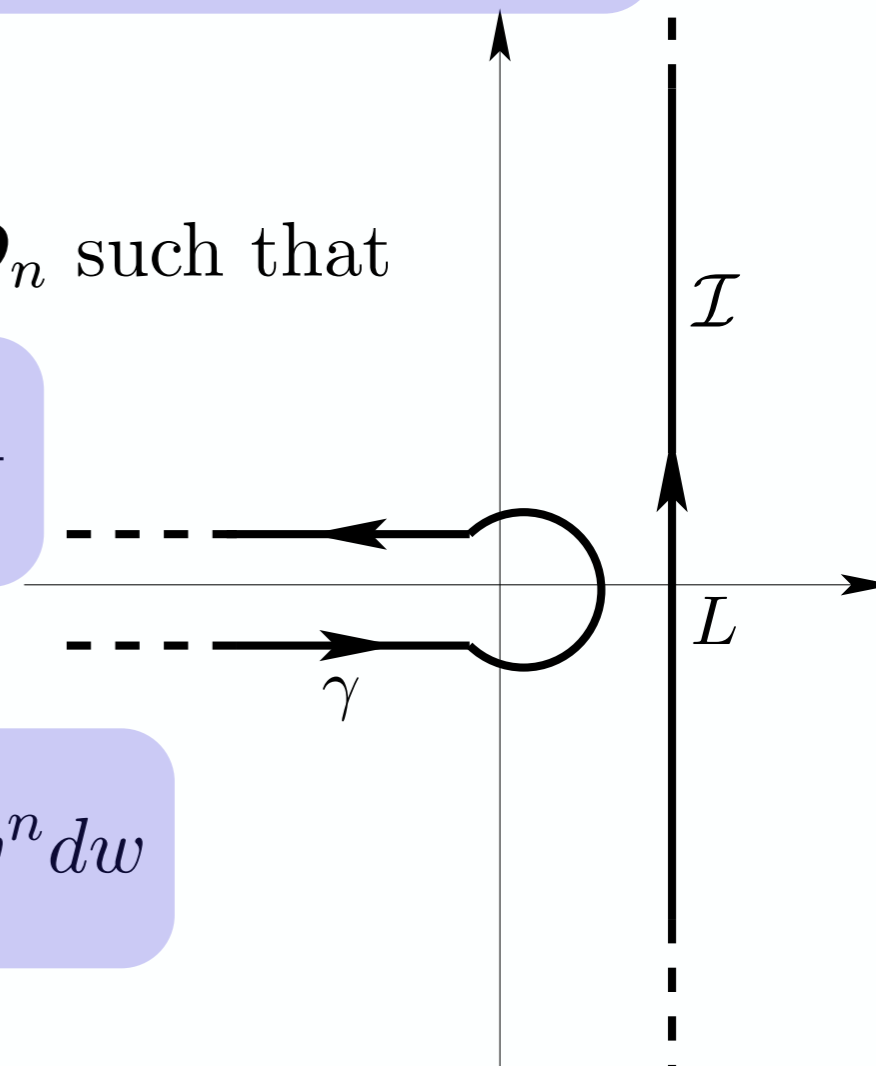


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$$P_n(x)e^{Ax} = e^{x^2} \int_{\mathcal{I}} w^J D_n w^{-J} e^{w^2 - 2xw} w^n dw$$



CASE $N=2$

For the family of MOP $(P_n)_n$ with respect to $(N = 2)$

$$\mathbf{W}(x) = e^{-x^2} \begin{pmatrix} 1 + x^2\nu^2 & \nu x \\ \nu x & 1 \end{pmatrix}, \quad x \in \mathbb{R}$$

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$$P_n(x) \begin{pmatrix} 1 & \nu x \\ 0 & 1 \end{pmatrix} = \frac{n!}{2^{n+1}\pi i} \oint_{\gamma} \begin{pmatrix} 1 & \frac{(n+1)\nu}{2z} \\ -\frac{z\nu}{\gamma_n^2} & \frac{1}{\gamma_n^2} \end{pmatrix} e^{-z^2+2zx} \frac{dz}{z^{n+1}}$$

$$P_n(x) \begin{pmatrix} 1 & \nu x \\ 0 & 1 \end{pmatrix} = \frac{e^{x^2}}{i\sqrt{\pi}} \int_{\mathcal{I}} \begin{pmatrix} 1 & w\nu \\ -\frac{n\nu}{2w\gamma_n^2} & \frac{1}{\gamma_n^2} \end{pmatrix} e^{w^2-2xw} w^n dw$$

CHRISTOFFEL-DARBOUX KERNEL (N=2)

The CD kernel

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$$\frac{2}{(2\pi i)^2} e^{\frac{x^2 - y^2}{2}} \int_{\mathcal{I}} dw \oint_{\gamma} dz z^{J_2} \mathbf{B}_n z^{-J_2} w^{J_2} \mathbf{B}_n^{-1} w^{-J_2} \frac{e^{w^2 - 2xw - z^2 + 2zy + n \log(w/z)}}{w - z}$$

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In other words

$$\frac{2}{(2\pi i)^2} e^{\frac{x^2-y^2}{2}} \int_{\mathcal{I}} dw \oint_{\gamma} dz \begin{pmatrix} \frac{z(\gamma_n^2-1)+w}{w\gamma_n^2} & \frac{\nu(w-z)}{\gamma_n^2} \\ \frac{n\nu(z-w)}{2\gamma_n^2} & \frac{w(\gamma_n^2-1)+z}{z\gamma_n^2} \end{pmatrix} \frac{e^{w^2-2xw-z^2+2zy+n \log(w/z)}}{w-z}$$

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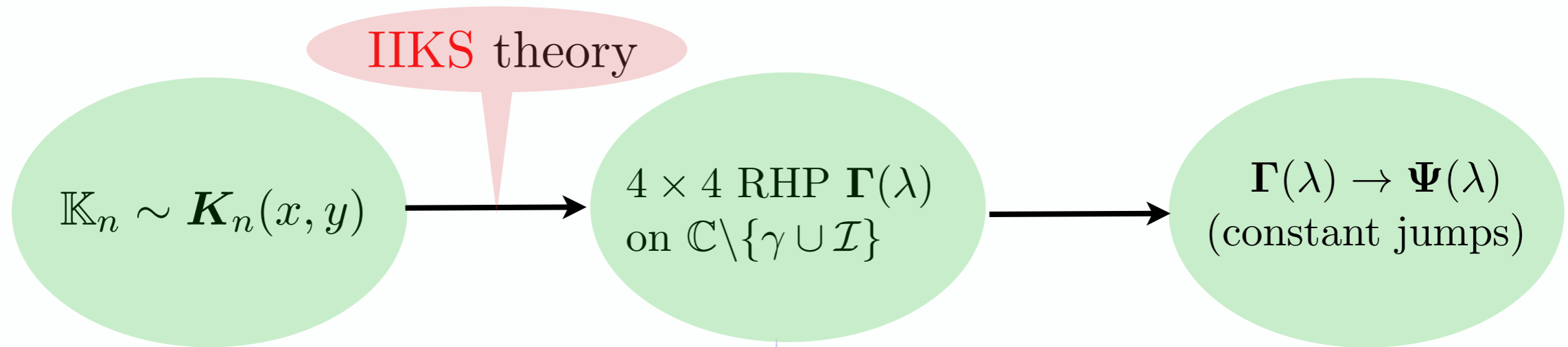
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$\mathbf{G}(\lambda)$ depends on $\lambda^{J_2} \mathbf{B}_n \lambda^{-J_2}$

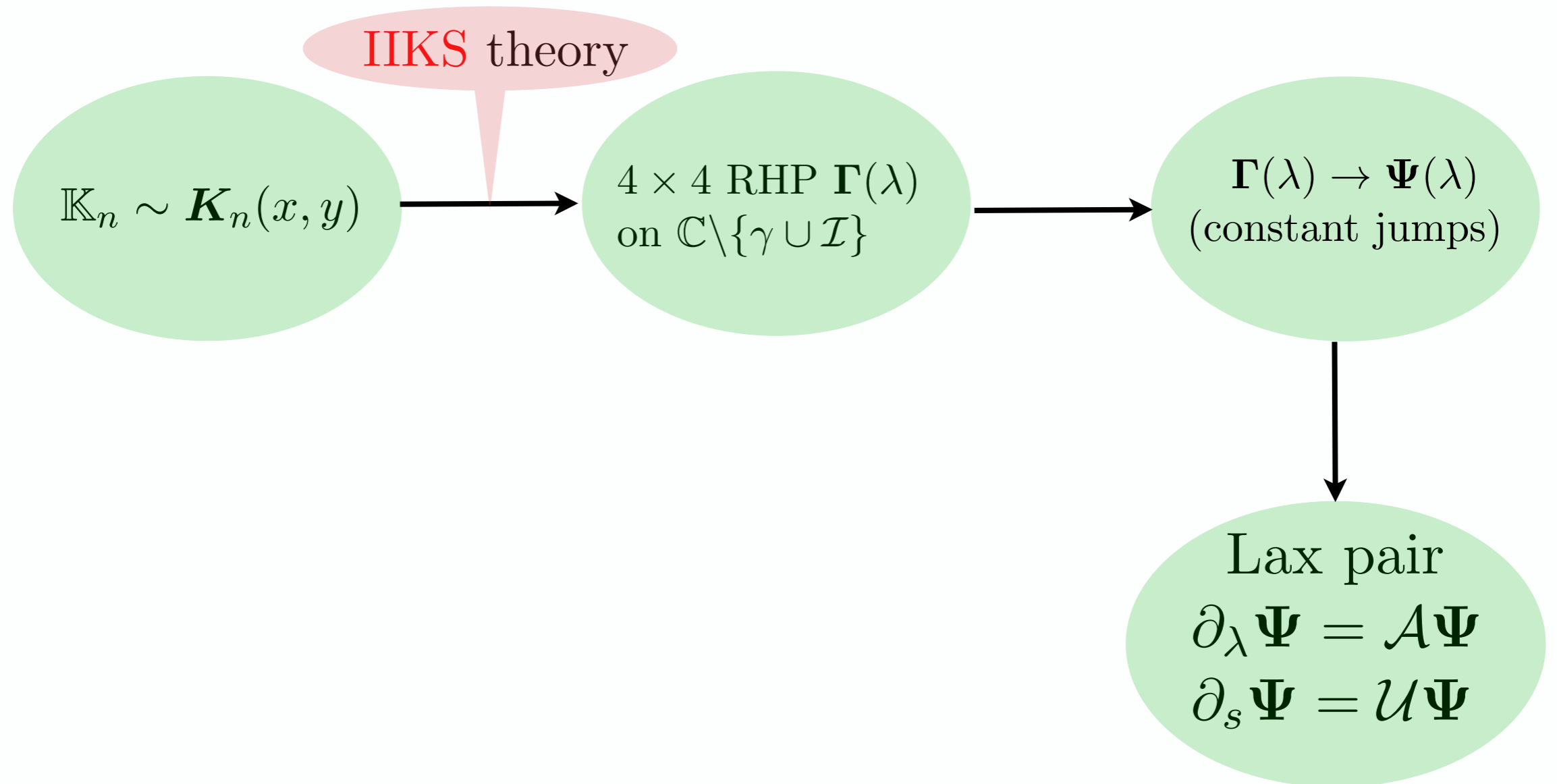
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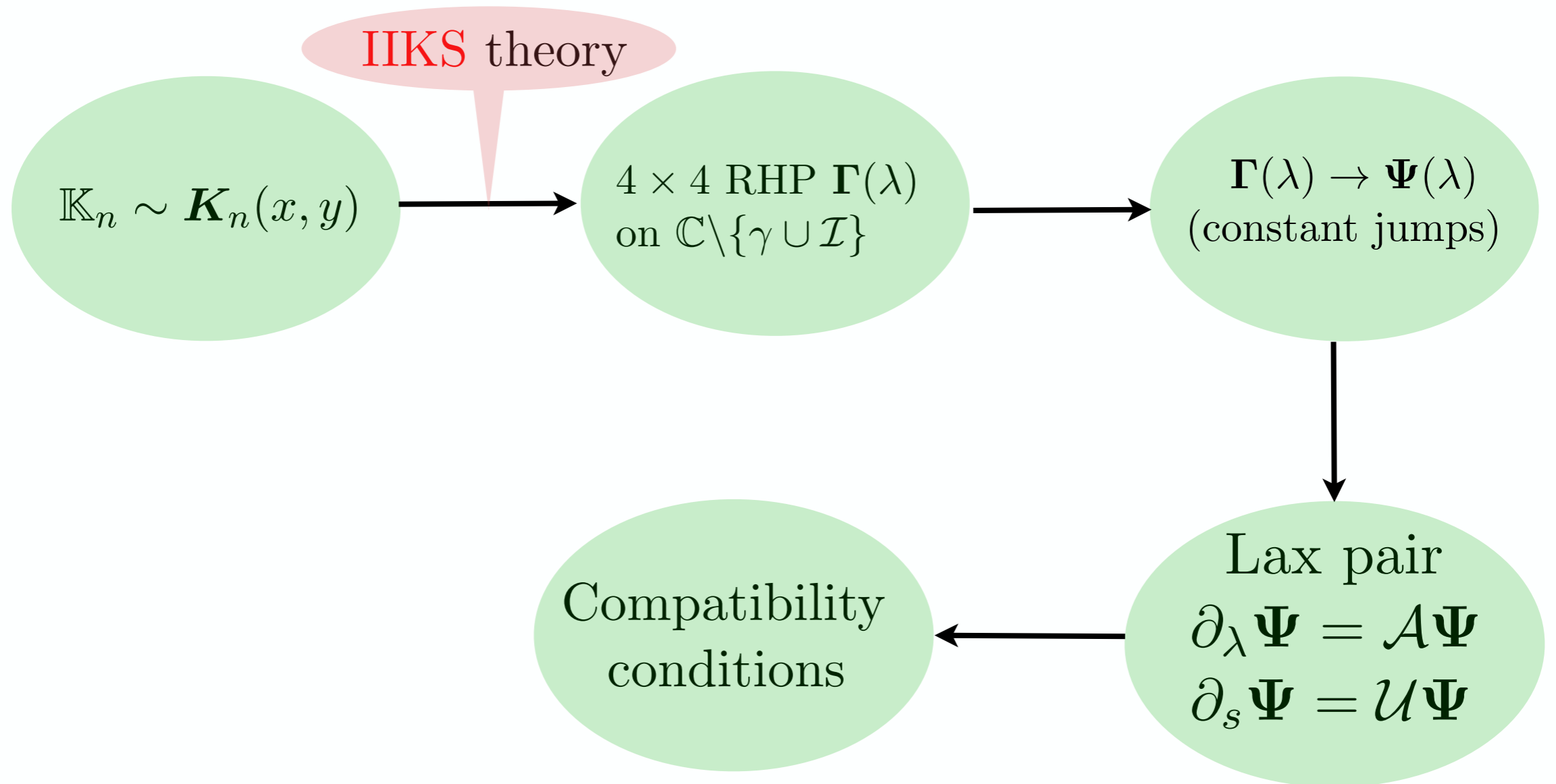
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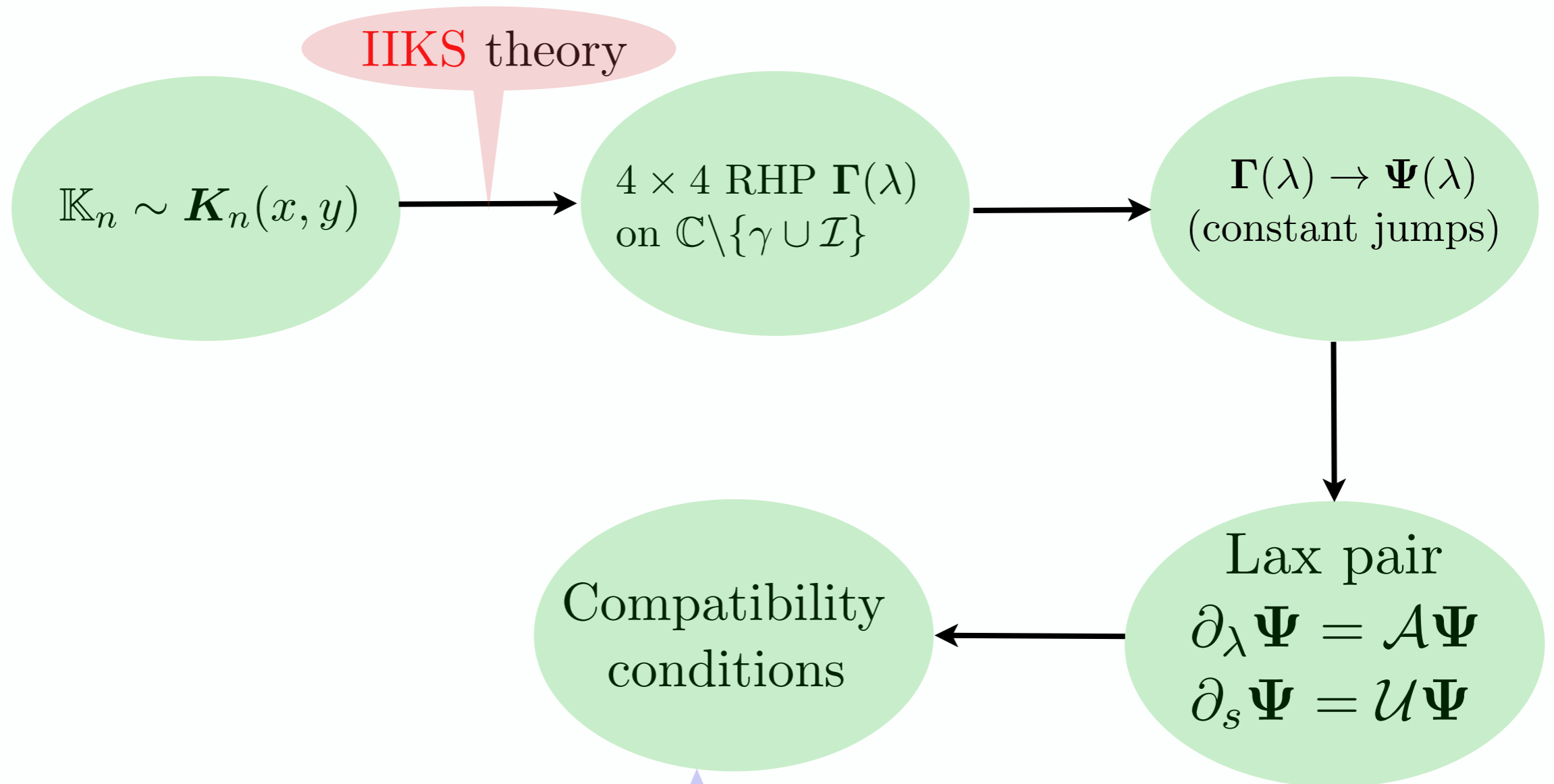
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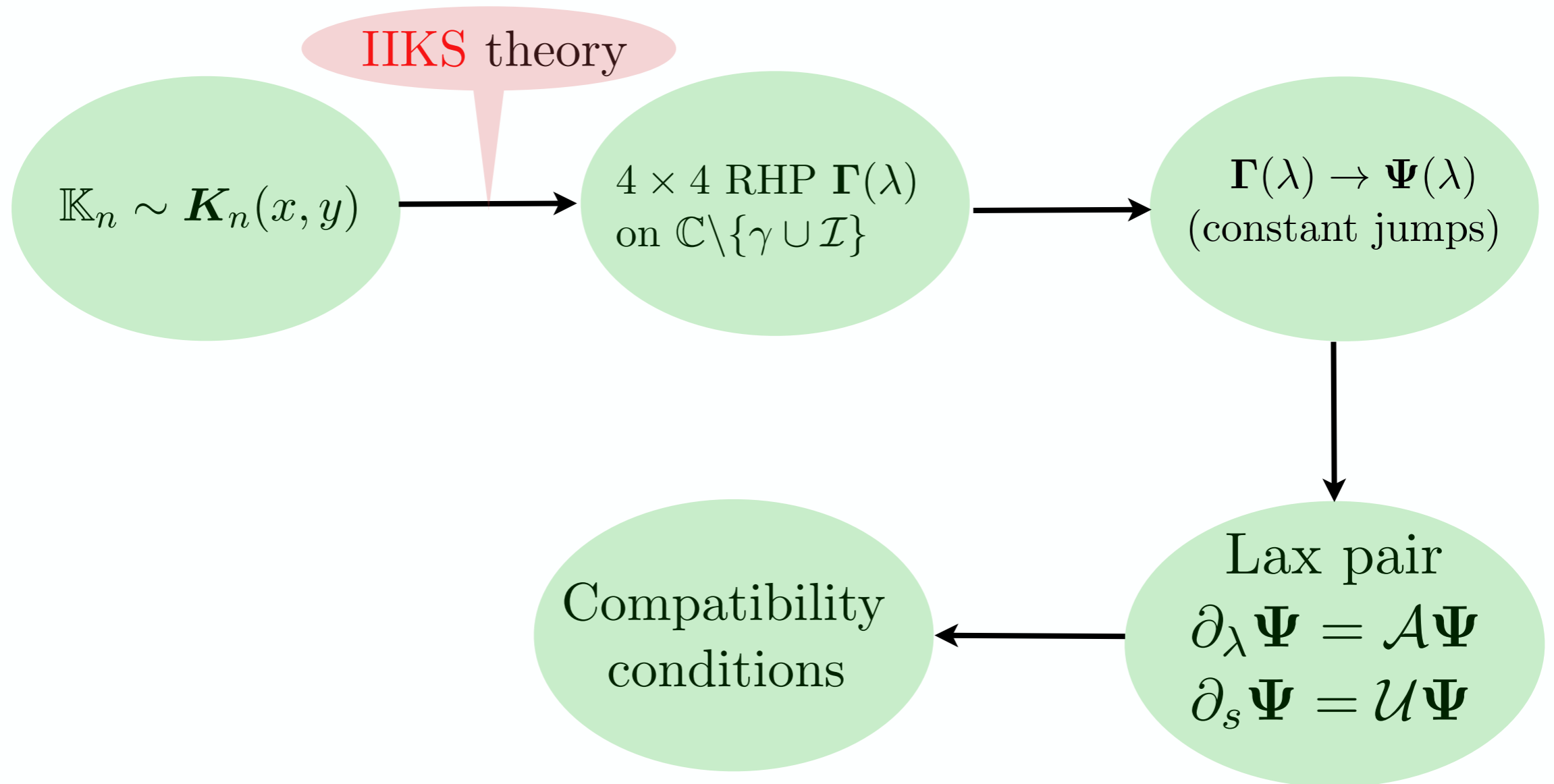


$$\begin{cases} u' &= -u^2 + 2su + 4z - 2nI_2 + V_A \\ z'' &= 2u'z + 2uz' - 2sz' \end{cases}$$

where $V_A = 2[\mathbf{J}_2, \mathbf{y}]\mathbf{y}^{-1}$ ($[\mathbf{x}, \mathbf{y}] = \mathbf{x}\mathbf{y} - \mathbf{y}\mathbf{x}$) and

$$z = -(\Gamma_1)'_{11}, \quad \mathbf{y} = -2(\Gamma_1)_{12}, \quad \mathbf{u} = (\Gamma_1)'_{12}(\Gamma_1)_{12}^{-1} + 2sI_2$$

NON-COMMUTATIVE PAINLEVÉ IV



Non-commutative version of the derived Painlevé IV equation

$$\begin{aligned} & \mathbf{u}''' + [\mathbf{u}'', \mathbf{u}] - 4(n + 1 + s^2)\mathbf{u}' - 2(\{\mathbf{u}', \mathbf{u}^2\} + \mathbf{u}\mathbf{u}'\mathbf{u}) \\ & + 6s\{\mathbf{u}', \mathbf{u}\} + 4\mathbf{u}(\mathbf{u} - s\mathbf{I}_2) + (\mathbf{V}'_A - 2(\mathbf{u}\mathbf{V}_A))' + 2s\mathbf{V}'_A = \mathbf{0} \end{aligned}$$

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If we assume that **all the variables commute**, we get the equation

$$u''' - 4u' - 6u^2u' + 12u'u - 4nu' + 4u^2 - 4su - 4s^2u' = 0$$

and this equation is the derivative of the Painlevé IV equation

$$u'' = \frac{(u')^2}{2u} + \frac{3}{2}u^3 - 4su^2 + 2(s^2 + 1 + n)u - \frac{2n^2}{u}$$

FINAL REMARKS

- It is well known that

$$\lim_{n \rightarrow \infty} \frac{1}{\sqrt{2n^{1/6}}} K_n^{\text{Hermite}} \left(\sqrt{2n} + \frac{x}{\sqrt{2n^{1/6}}}, \sqrt{2n} + \frac{y}{\sqrt{2n^{1/6}}} \right) = K_{\text{Ai}}(x, y),$$

In our case we have that

$$\lim_{n \rightarrow \infty} \frac{1}{\sqrt{2n^{1/6}}} \mathbf{K}_n \left(\sqrt{2n} + \frac{x}{\sqrt{2n^{1/6}}}, \sqrt{2n} + \frac{y}{\sqrt{2n^{1/6}}} \right) = K_{\text{Ai}}(x, y) \mathbf{I}_N.$$

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- M. D. de la Iglesia and M. Cafasso, *Non-commutative Painlevé equations and Hermite-type matrix orthogonal polynomials*, accepted in *Communications in Mathematical Physics*.



THANK YOU

